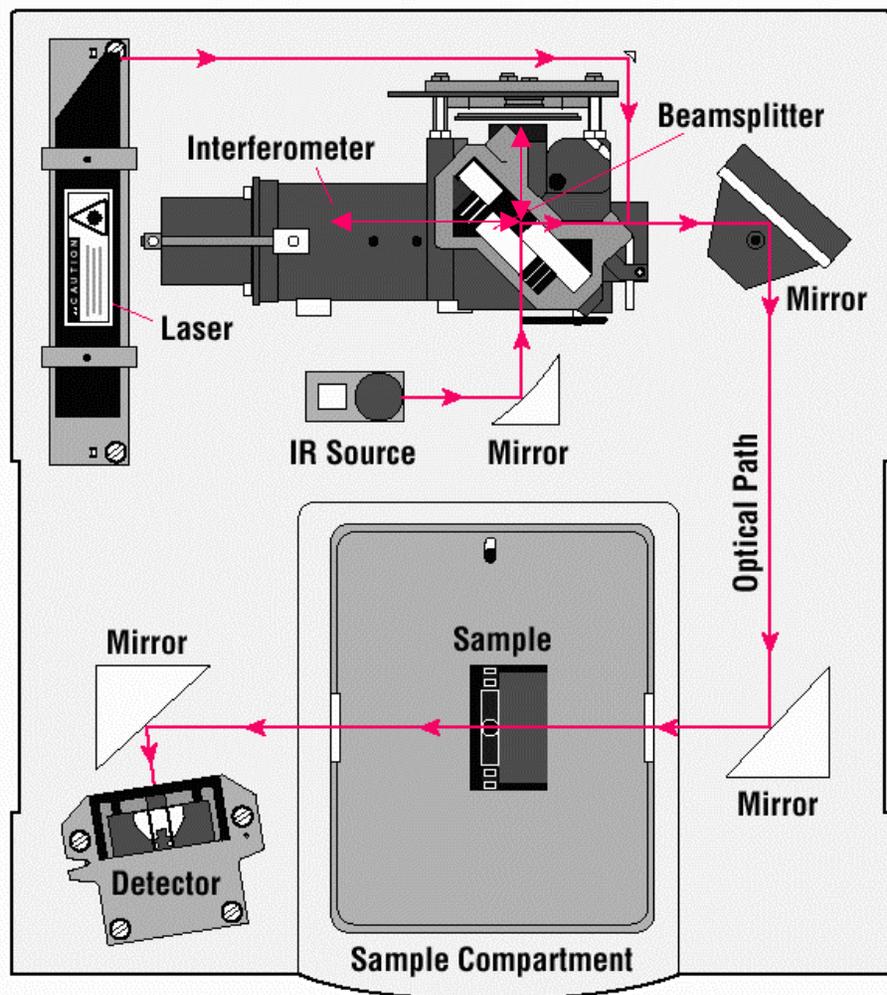




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How an FTIR Spectrometer Works

A Simple Spectrometer Layout



Pathlength difference = x

The intensity detected of two plane waves:

$$I = |\vec{E}|^2 = |E_1|^2 + |E_2|^2 + 2\vec{E}_1 \cdot \vec{E}_2 \cos(\mathbf{q})$$

Normal incidence, $\theta = kx$, can simplify to:

$$I(x) = 2[1 + \cos(kx)]$$

For non-monochromatic light:

$$\begin{aligned} I(x) &= \int_0^{\infty} [1 + \cos(kx)] G(k) dk \\ &= \int_0^{\infty} G(k) dk + \int_0^{\infty} G(k) \frac{e^{ikx} + e^{-ikx}}{2} dk \\ &= \frac{1}{2} I(0) + \frac{1}{2} \int_{-\infty}^{\infty} G(k) e^{ikx} dk \end{aligned}$$



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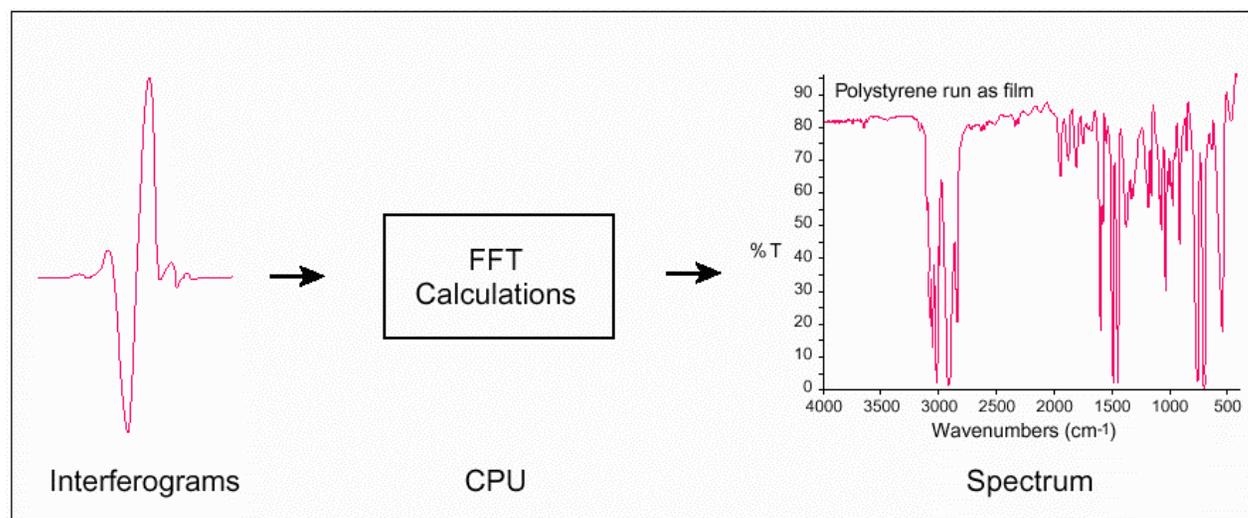
FTIR Math Continued

We can rewrite this to something more familiar:

$$W(x) \equiv \frac{2I(x) - I(0)}{\sqrt{2p}} = \frac{1}{\sqrt{2p}} \int_{-\infty}^{\infty} G(k) e^{ikx} dk$$

A Fourier Transform!

The detected intensity as a function of moving mirror position, $I(x)$, can therefore be converted into $G(k)$, the intensity spectrum as a function of frequency by a simple Fourier transform.





FTIR Spectrometers

In practice one cannot measure from $-\infty$ to ∞ . The resolution of a measurement is simply given by how far in x you measure.

$$\text{resolution} \propto \frac{1}{2px_{\max}}$$

Rapid-Scan measurements:

- Sweep mirror quickly, average many interferograms
 - Very fast & easy
 - Not high resolution
 - Not for quickly changing signals or very low signal

Step-Scan measurements:

- Step to each x position, then measure (long average, or triggered time series). Can have very long path length.
 - Excellent for fast time resolution, low signals (lock-in)
 - Harder to run stably.